Original Research Paper

Technical Stability of Very Slender Rectangular Columns Compressed by Ball-And-Socket Joints without Friction

Krzysztof Murawski

Independent Researcher, Poland

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Email: k.murawski@interia.pl

Abstract: Stability of load-bearing members, as known, is a challenging issue and several tools are available for designers. Disregarding the material properties in use, the avoidance of possible stability troubles is a mandatory and challenging step of the overall design process. In this study, a theoretical model is presented for the of loss of stability in elastic states of very slender rectangular shell elements axially compressed through balland-socket joints without friction. According to this theory, as shown, a loss of carrying capacity of very slender columns in elastic states occurs when the line of force leaves a critical transverse cross-section. The critical transverse cross-section, moreover, progressively moves because of the superposition of bending and pure compression. The theory allows to determine the governing differential equations of curved central lines and their slopes, as well as the critical stresses of columns, in the form of surface function in dependence on slenderness ratios and cross-sectional areas. The graphs for the elastic deflected central line y(x), slope dy/dx, dependence $y_{L/2}(P)$, stresses and strains for a rectangular column made of steel and compressed by ball-and-socket joints without friction, as well as the corresponding surface graphs of critical stress, are presented in this study. The obtained surface graphs of critical stresses are then discussed and compared with Euler's formulation.

Keywords: Stability, Elastic States, Slenderness, Rectangular Columns, Steel

Introduction

Shell structures, in the field of building construction and engineering in general, represent a thin, curved plate structure that is optimally shaped to transmit applied forces by compressive, tensile and shear stresses that act in the plane of the same surface. When we consider the application of shell elements in load-bearing structures, accordingly, the first issue to analyze is their load capacity to sustain axial loads, i.e., their stability and susceptibility to potential buckling collapse mechanisms that could compromise the structure and occupant safety. In the case of very slender columns, this refers to the problem of stability in elastic states.

The basic theory of slender rods losing stability in elastic states, as known, has been originally formulated by (Euler, 1744; 1759). He first introduced the concept of critical load P_{cr} and presented, according to his theory, the differential equation of an elastic deflected central line, that is:

$$EJ\frac{d^2y}{dx^2} = -P_{cr}y,\tag{1}$$

Where:

- E: Denotes the Young's modulus of elasticity of the column
- J: The moment of inertia of the cross-section area, while
- *y*: Represents the distance from the undeformed central line, from the *y*-axis

He had assumed a displacement y(x) of the column axis as a part of a sine curve:

$$y(x) = \sin\left(n \cdot \frac{\pi}{L}x\right),\tag{2}$$

Where:

n: Denotes the coefficient of the sine curve part and depends on boundaries (n = 1 for a pinned column, n



= 2 for a column fixed at one end, n = 1/2 for a column fixed by a pinned end and a vertically slide for the second end, n = 0.7 for a column with a pinned end and a pinned vertical slide at the second end)

L: The length of the bar

And he has obtained the formula for a critical stress σ_{cr} of an axially compressed column by force as follows:

$$\sigma_{cr}^{Euler} = \left(\frac{\pi}{n \cdot \lambda}\right)^2 E.$$
(3)

The problem of the stability was later searched and analyzed further with a focus on many other relevant aspects for various engineering applications, like buckling of columns composed of various constructional materials (Gupta et al., 2001; Oleiwi et al., 2014; Schnabl et al., 2013; Kalamar et al., 2016; Wahrhaftig et al., 2016; Thermou et al., 2018; Zhou et al., 2017; Ye et al., 2018; Osmani and Meftah, 2018; Wang et al., 2018; Naderpour et al., 2019; Qi et al., 2019; Goroshko et al., 2020), as well as beams (Ascione and Grimaldi, 1983; Di Sarno and Manfredi, 2012; Li et al., 2015; Özbaşaran et al., 2015; Monsalve-Cano and Aristizábal-Ochoa, 2016; Nguyen et al., 2018; Belaid et al., 2018) or nanobeams (Rahmani et al., 2017; Mohammadi et al., 2019).

Single-layer graphene sheets have been examined in (Genoese *et al.*, 2019), while polymer-confined concrete columns have been discussed in (Liang *et al.*, 2012) and hyperelastic tubes are analyzed by (Liu, 2018). Many other research contributions have been then related to the stability issues of a multitude of loadbearing systems and members can be found in the literature, including plates (Sabouri-Ghomi *et al.*, 2008; Rao and Ra, 2009; Xu *et al.*, 2013; Moradi-Dastjerdi and Malek-Mohammadi, 2017; Riahi *et al.*, 2018; Vu *et al.*, 2019) and nanoplates (Malikan *et al.*, 2018), bracing systems (Rahnavard *et al.*, 2018), tubes (Nouri *et al.*, 2015; Mozafari *et al.*, 2018; Naveed *et al.*, 2017; Sadath *et al.*, 2017; Sun *et al.*, 2018), frames (Marante *et al.*, 2012; Slimani *et al.*, 2018), pipes (Lolov and Lilkova-Markova, 2005; Melissianos *et al.*, 2017; Moustabchir *et al.*, 2018; Psyrras *et al.*, 2019), or Functionally Graded Material (FGM) structures (Moita *et al.*, 2018; Kiss, 2019; Singh and Harsha, 2019), etc.

Stability Analysis

This research study herein discussed assumes as its own method of analysis for the column stability in elastic states that the state of stresses and strains in a critical transverse section (after losing stability and before losing its carrying capacity) appears as a superposition result of pure compression phenomena and bending effects (Fig. 1 and also (Murawski and Kłos, 2007; Murawski, 1992; 2011; 2018)). The method is thus developed on a basic simplification for very slender columns, that is the beginning of the load-carrying capacity lost (i.e., the maximum achieved value of force, on a force *P*-shortening ΔL graph) in elastic states follows the exceedance of the force line from a critical transverse section.



Fig. 1: Stress and strain analysis in the critical transverse section for a rectangular column axially compressed by force through balland- socket joints, after the loss of stability and before the loss of carrying capacity

The fibre extension ε at the distance *y* from the undeformed central line and the corresponding stress σ_n are in fact given by (Fig. 1):

$$\varepsilon = \frac{\pm (\rho + y) \cdot \theta - \rho \cdot \theta}{\rho \cdot \theta} = \pm \frac{y}{\rho},\tag{4}$$

$$\sigma_n = \sigma_g - \sigma_c = \pm \frac{y}{\rho} \cdot E - \frac{P}{A},$$
(5)

where, E is defined in Equation (1), while:

- ρ : Denotes the radius of the curved central line
- θ . The angle of the central line slope, in relation to the force line
- P: The imposed axial force
- A: The area of the critical cross-section
- σ_g : Represents the bending stress
- σ_c : The compressive stress
- σ_n : The normal stress

Due to the force equilibrium, it is:

$$dP = \left(\sigma_g + \sigma_c\right) \cdot dA = \left(\pm \frac{y}{\rho} \cdot E - \frac{P}{A}\right) \cdot dA.$$
 (6)

In this study, unlike Euler, it is moreover assumed that the displacement of the column axis is determined by taking into account the conditions of the small deformation theory, as a function of the curvature radius, that is:

$$\frac{1}{\rho} \cong \frac{d^2 y}{dx^2},\tag{7}$$

and the reference boundary conditions:

$$dy/dx = 0 \text{ when } x = L/2, \tag{8}$$

and:

$$y = 0 \quad \text{when} \quad x = 0, \tag{9}$$

The governing differential equations for the curved central line d^2y/dx^2 and its slope dy/dx, as well as the equation of the central line y(x) for a column axially compressed by force through ball-and-socket joints without friction are like Equation (59), (60), (61) in (Murawski, 2011), or Equation (2.7), (2.8), (2.9) in (Murawski, 2018). From the assumption that the losing of carrying capacity follows when the force line exceeds the critical transverse section ($y_x = L/2 = y_{cr}$), moreover, the critical stress σ_{cr} for axially compressed columns under force through ball-and-socket joints is like Equation (65) in (Murawski, 2011), or Equation (2.13) in (Murawski, 2018).

Theoretical Example

As a theoretical example, the graphs of the elastic line y(x), its slope dy/dx and the dependence $y_{L/2}(P)$ for a rectangular column made of steel are presented. In doing so, a steel column agreeing with Fig. 1 is taken into account.

Moreover, the nominal dimensions of the examined column are set in a = 20 mm, b = 28 mm and t = 1 mm for the resisting cross-section, while L = 2500 mm is the selected span (and $\lambda = 314.8$ the corresponding slenderness) and the column is compressed by ball-and-socket joints. The so-collected results are presented in Fig. 2.

Under the assumption that the loss of carrying capacity occurs when the force line leaves the critical transverse section $(y_{x=L/2} = y_{cr} = a/2+t)$, the dependency of various geometrical parameters on the mechanical performance of the selected column is highlighted in Figs. 3a-3d.



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Fig. 2: Evolution of (a) elastic line y(x), (b) elastic line slope dy/dx and (c) function $y_L(P)$ for a rectangular steel column (dimensions: a = 20 mm, b = 28 mm, t = 1 mm, L = 2500 mm) compressed by ball-and-socket joints



(a)

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(b)



 $\sigma_{cr}(\zeta, \chi, \mu = a/t = 20, t = 1)$

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Fig. 3: Surface functions of critical stress for rectangular shaped shells made of steel and axially compressed by force through balland-socket joints, according to the technical stability theory



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Fig. 4: Surface functions of critical stress for axially compressed (by force through ball-and-socket joints) rectangular shaped shells made of steel, according to the technical stability theory

The focus in Fig. 3 is given to:

- a) The critical stress that is shown in Fig. 3a, as a function of the cross sectional area A and slenderness ratio λ at $a/b = \cos t$ and $t/b = \cos t$
- b) The function shown in Fig. 3b is the same of Fig. 3a, but the parameters $\mu = a/t$, $\zeta = b/t$ and $\chi = a/L$ are used under the condition that $t = \cos t$
- c) The same function is shown in Fig. 3c, with the exception that t = cost with the parameter μ constant, but with χ and μ modify
- d) Finally, the function is shown in Fig. 3d with constant parameters μ = a/t and ζ = b/t, but variable χ and t

The critical stress can be described also by the empirical formulas presented in Fig. 4.

For thin walled columns (t > 0), the corresponding formulas for the critical stress analysis are presented in Fig. 5. Again, the sensitivity of critical stress is shown as a function of:

a) the cross sectional area *A* and slenderness ratio λ at $a/b = \cos t$

b) slenderness ratio λ and L/a = var at a/b = cost

For comparative purposes, the Euler's formulas of critical stresses for rectangular columns in the dependence on μ , χ , ζ and *t* are then presented in Fig. 6. More in detail, the analysis is proposed as a function of:

- a) The slenderness ratio λ
- b) The parameters $\mu = a/t = \text{var}, \ \chi = a/L = \text{var}, \ \zeta = b/t = \text{cost}$
- c) The parameters $\zeta = b/t = var$, $\chi = a/L = var$, $\mu = a/t = cost$
- d) And finally the parameters t = var, L = var, a = cost, with b = cost

In the case of the reference rectangular column (a = 20 mm, b = 28 mm, t = 1 mm, L = 2500 mm) the critical stress according to Euler's theory is calculated in $\sigma_{cr}^{Euler}(\lambda) = 16.592 \text{ MPa}$. The reference value of Young's modulus is set in E = 166600 MPa, according to past stability studies on thin walled columns (Murawski, 2011). In this regard, the critical stress for the same column, according to the technical stability theory, is predicted in $\sigma_{cr}^{balls}_{rectangular}(\lambda, A, a/b, t/b) = 3.562 \text{ MPa}$.





Fig. 5: Surface functions of the critical stress of axially compressed (by force through ball-and-socket joints) rectangular shaped shells made of steel, according to the technical stability theory





(b)

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Fig. 6: Surface function of the critical stress based on Euler's formulation of rectangular shaped shells made of steel and axially compressed by force through ball-and-socket joints



Fig. 7: Surface function of the limited slenderness ratio $\lambda_{el\cdot lt}$ based on the technical stability formula of rectangular shaped shells made of steel in dependence on (a) cross-sectional area *A*, ratio b/a = var and t/a = cost, (b) cross section area *A*, ratio t/a = var and b/a = cost



Fig. 8: Surface function of limited slenderness ratio λ_{el-ll} , based on the technical stability formulation, of thin walled rectangular columns made of steel, in dependence on *A* and a/b = var

Since the described phenomenon occurs in an elastic state for steel shells, it is thus expected that their slenderness ratio should be at least $\lambda > \lambda_{el-lt}$, according to Euler's theory.

The limited slenderness ratio λ_{el-lt} is then calculated according to the technical stability theory in Fig. 7. For thin walled columns the formula for the limited slenderness ratio λ_{el-lt} thin_walled can be also seen in Fig. 8.

Stress and Strain Analysis

According to the technical stability theory, the state of stresses and strains in the shell belonging to the critical transverse section of a rectangular shaped column compressed by ball-and-socket joints, after the loss of stability, is as the result of a superposition of bending and pure compression (Fig. 1) and thus:

$$\varepsilon_{n} \cdot E = \sigma_{n} = \sigma_{g} + \sigma_{c} = \pm \frac{y}{\rho} \cdot E - \frac{P}{A}$$

$$= E \cdot \left(\pm \frac{y}{\rho} - \frac{\Delta L}{L} \right) = E \cdot \left(\pm \varepsilon_{g} - \varepsilon_{c} \right) \Longrightarrow$$
(10)

with:

$$\Longrightarrow \mathcal{E}_n = \pm \mathcal{E}_g - \mathcal{E}_c. \tag{11}$$

It results that the shell normal stress σ_n is like Equation (71) in (Murawski, 2011), or Equation (2.19) from (Murawski, 2018):

The graphs of the functions of the normal shell stress σ_n , normal strain ε_n , transverse stress σ_y and transverse strain ε_y in the selected rectangular column (a = 20 mm, b = 28 mm, t = 1 mm, L = 2500 mm), axially compressed by force through ball-and-socket joints, are presented in Figs. 9, 11 and 13.

The stresses and strains are proposed in dependence of x and y, where $x = 0 \div L$ and y = -a/2, -a/4, -a/20, 0.0, a/20, a/4, a/2, respectively.

The graphs of the functions of the normal stress $\sigma_{n_x=L/2}(y)$, normal strain $\varepsilon_{n_x=L/2}(y)$, transverse stress $\sigma_{y_x=L/2}(y)$ and transverse strain $\varepsilon_{y_x=L/2}(y)$ in the critical cross-section (x = L/2), depending on y and x, are presented in Figs. 10, 12 and 14 for load values set in P= 0.1, 0.5, or 1.0 $P_{\rm cr}$.

The graphs of the function of the coordinate $y_{\sigma_{n=0,x}}$ x=L/2(P) of a zero normal stress $\sigma_{n_x}=L/2(y) = 0$ in the critical cross-section (x = L/2) in dependence on the load *P* are presented in Fig. 15a and the graphs of the function of the deformation $\Delta L_{x=L/2,y=a/2}(P)$ in the critical (middle) cross-section (x = L/2) of the rectangular shaped steel shell on two opposite generating lines ($y = \pm a/2$) in

dependence on the load *P* is presented in Fig. 15b. As also previously described, the values of Young's modulus E = 166600 MPa and Poisson's ratio v = 0.3 were taken into account for calculations.



Fig. 9: Values of (a) stress $\sigma_n(x,y)$, strain $\varepsilon_n(x,y)$, deformation $\Delta L(x,y)$ and (b) stress $\sigma_y(x,y)$, strain $\varepsilon_y(x,y)$, deformation $\Delta a(x,y)$ where $x = 0 \div L$ and y = -a/2, -a/4, -a/20, 0.0, a/20, a/4, a/2 respectively, for $P = 0.1 P_{cr}$ in a rectangular steel shell (dimensions: a = 20 mm, b = 28 mm, t = 1 mm, L = 2500 mm) compressed by ball-and-socket joints

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Fig. 10: Values of (a) normal stress-strain-deformation $\sigma_n, \varepsilon_n, \Delta L(y, x = L/2)$ and (b) transverse stress-strain-deformation $\sigma_y, \varepsilon_y, \Delta 2t(y, x = L/2)$ for $P = 0.1 P_{cr}$ in a rectangular steel shell (dimensions: a = 20 mm, b = 28 mm, t = 1 mm, L = 2500 mm) compressed by ball-and-socket joints



Fig. 11: Values of (a) stress $\sigma_n(x,y)$, strain $\varepsilon_n(x,y)$, deformation $\Delta L(x,y)$ and (b) stress $\sigma_y(x,y)$, strain $\varepsilon_y(x,y)$, deformation $\Delta a(x,y)$ (where $x = 0 \div L$ and y = -a/2, -a/4, -a/20, 0.0, a/20, a/4, a/2 respectively) for an imposed load value $P = 0.5 P_{cr}$, in a rectangular steel shell (dimensions: a = 20 mm, b = 28 mm, t = 1 mm, L = 2500 mm) compressed by ball-and-socket joints



Fig. 12: Values of (a) normal stress-strain-deformation $\sigma_n, \varepsilon_n, \Delta L(y, x = L/2)$ and (b) transverse stress-strain-deformation $\sigma_y, \varepsilon_y, \Delta 2t(y, x = L/2)$ for $P = 0.5 P_{cr}$, in a rectangular shell made of steel (dimensions: a = 20 mm, b = 28 mm, t = 1 mm, L = 2500 mm) compressed by ball-and-socket joints

Discussion of Results

The main conclusion which can be drawn out from Fig. 3a is that in order to determine the critical stresses $\sigma_{cr}^{balls}_{rectangular}$ for general rectangular shaped shell elements, the slenderness ratio λ is not sufficient. The critical stress estimates should be indeed related also to

the cross-section area A and to values of the section ratios a/b and t/b.

For a given slenderness ratio λ like in original Euler's formulation) it is possible to obtain many different values of the critical stress $\sigma_{cr}^{balls}_{rectangular}$ in dependence of the cross-section area *A* and the ratios a/b and t/b.



Fig. 13: Values of (a) stress $\sigma_n(x,y)$, strain $\varepsilon_n(x,y)$, deformation $\Delta L(x,y)$ and (b) stress $\sigma_y(x,y)$, strain $\varepsilon_y(x,y)$, deformation $\Delta a(x,y)$ for $P = 0.5 P_{cr}$, in a rectangular shell made of steel (dimensions: a = 20 mm, b = 28 mm, t = 1 mm, L = 2500 mm) compressed by ball-and-socket joints (where $x = 0 \div L$ and y = -a/2, -a/4, -a/20, 0.0, a/20, a/4, a/2 respectively).



Fig. 14: Values of (a) normal stresses-strains σ_n , ε_n , $\Delta L(y, x = L/2)$ and (b) transverse stress-strains σ_y , ε_y , $\Delta 2t(y, x = L/2)$ for P = 1.0 P_{cr} in a rectangular steel shell (dimensions: a = 20 mm, b = 28 mm, t = 1 mm, L = 2500 mm) compressed by ball-and-socket joints

The theoretical results, moreover, show that the simplifications used in the technical stability theory like taking under consideration only normal stresses, the assumption of flat cross-section planes and small slope of central line, involves certain limitations on its application in the engineering practice. The first of those limitations, for rectangular shaped shells, is that the cross-section area *A* must be bigger than the unit.

Otherwise, the function approaches asymptotically to the vertical flat plane A = 1 (Fig. 3a).

The surface function in Fig. 3a with a/b = 0.714, t/b = 0.0357 equals Euler's function $\sigma_{cr}^{Euler}(\lambda)$ for A = 2.6327.



Fig. 15: Values of (a) coordinate $y_{\sigma n=0, x=L}(P)$ of zero normal stress $\sigma_n = 0$ and (b) deformation $\Delta L_{x=L,y=\pm a/2}(P)$ in the critical crosssection (x = L/2), in dependence of *P* for an axially compressed rectangular shaped shell made of steel (dimensions: a = 20 mm, b = 28 mm, t = 1 mm, L = 2500 mm).

Because of this reason, it is found that: $\sigma_{cr}^{Euler}(\lambda) > \sigma_{cr}^{balls_{rectangular}}(\lambda, A, a/b = 0.714, t/b = 0.0357)$ for A < 2.6327 and $\sigma_{cr}^{Euler}(\lambda) < \sigma_{cr}^{balls_{rectangular}}(\lambda, A, a/b = 0.714, t/b = 0.0357)$ for A < 2.6327.

Some conclusions that can be thus drawn from the analysis of the surface function for $\sigma_{cr}^{balls}_{rectangular}(\mu, \chi, \zeta = 28, t = 1)$ that was shown in Fig. 3b.

With the increase of μ and the decrease of χ , in particular, it is observed that the stress value $\sigma_{cr}^{balls}_{rectangular}(\mu,\chi,\zeta=28, t=1)$ also increases and this is in agreement with the original Euler's formulation, as also shown in Fig. 6b. The graph in Fig. 3c, moreover, shows that with the increase of both the ζ and χ parameters, the measured stress value $\sigma_{cr}^{balls}_{rectangular}$

 $(\mu, \chi, \zeta = 20, t = 1)$ decreases at constant values of the ratio $\mu = a/t$ and thickness *t*.

The graph in Fig. 3d, finally, shows that with the increase of the ratio b/a and the increase of the ratio t/a, the stress $\sigma_{cr}^{balls}_{rectangular}(t, \chi, \mu = 20, \zeta = 28)$ decreases at constant values of the ratio $\chi = L/t$ and the thickness *t*.

Figure 4b shows that with the increase of the ratio b/a and with the increase of ratio t/a, the stress $\sigma_{cr}^{balls}_{rectangular}(b/a, t/a, \lambda = 314.8, a = 20)$ decreases at constant values of slenderness ratio λ and a.

From the analysis of the surface function $\sigma_{cr}^{balls_thin-}$ walled rectangular (λ , A, a/b=0.714) shown in Fig. 5a and of the surface function $\sigma_{cr}^{balls_thin-walled}$ rectangular (λ , A, a/b = 0.714) that is proposed in Fig. 5b, it is then possible to observe that - in case of thin walled columns - with the increase of the slenderness ratio λ and cross-section area A at constant values of the ratios a/b, the measured stress value $\sigma_{cr}^{balls_thin-walled}$ rectangular(λ , A, a/b = 0.714) also decreases.

The surface functions $\sigma_{cr}^{Euler's_balls}_{rectangular} = (\mu, \chi, \zeta)$ and $\sigma_{cr}^{Euler's_balls}_{rectangular} = (a, b, L, t)$, based on Euler's formulation, are not sensitive to the cross-section area A.

However, the Euler's formula does not distinguish the obtained results in dependence on the profile of the cross-section for the examined column.

This is indeed something that the technical stability formulation can capture.

On the graph in Fig. 7a, it is possible to see that the limited slenderness ratio $\lambda_{el\cdot lt}(A, b/a, t/a = 0.05)$, with the increase of the b/a ratio, minimally decreases, while strongly decreases at constant values of the ratio t/a with the increase of the cross-section area *A*. On the graph in Fig. 7b we see that the limited slenderness ratio $\lambda_{el\cdot lt}(A, t/a, b/a = 1.4)$ with the increase of the cross-section areas *A* decreases strongly at constant values of the ratio b/a.

On the graph in Fig. 8 we see that the limited slenderness ratio λ_{el-lt} ^{thin-walled} (A, t/a, b/a = 1.4) with the increase of the ratio a/b increases and with the increase of the cross-section area A strongly decreases.

On the graphs in Figs. 9a, 10a, 12a and 13a it is possible to see that the values of the stresses $\sigma_n(x,y)$, strains $\varepsilon_n(x,y)$ and deformations $\Delta L(x,y)$ for $P = 0.1 P_{cr}$ and 0.5 P_{cr} in the rectangular shell made of steel with dimensions: a = 20 mm, b = 28 mm, t = 1 mm, L = 2500mm compressed by ball-and-socket joints have minus sign, i.e., all fibers in parallel to the column axis direction are compressed. That means also that all fibers in perpendicular to the direction of the column axis direction are tensioned (Figs. 9b, 10b, 11b and 12b).

From the graphs in Figs. 13a and 14a it can be noticed that the values of stresses $\sigma_n(x,y)$, strains $\varepsilon_n(x,y)$ and deformations $\Delta L(x,y)$ in the rectangular shell (for $P = P_{cr}$) have partially negative sign. In this region, the fibers that are parallel to the column axis are compressed, while they are tensioned in the direction perpendicular to the column axis (Figs. 13b and 14b).

In the region where the stress, strain, deformation values have positive sign, the fibers in direction parallel to the column axis are tensioned while they are compressed in the direction perpendicular to the column axis direction (Figs. 13b and 14b).

The phenomena observed during past experimental tests (Murawski and Kłos, 2007) proved that the extensions of the two opposite generating lines are negative at the beginning of the phenomenon of the loss of stability. With the increase of time, however, one of them goes to positive values, as also shown on the graph in Fig. 16. This observation is in line with the technical stability theory (Fig. 15a) and this proves the accuracy of the proposed theory.



Fig. 16: (a) Extensions ΔL in time for a pine column ($870 \times \phi 12$ mm), as obtained from (b) the strain gauges placed on the two opposite generating lines (b) in the mid-span cross-section (Murawski and Kłos, 2007).

The graph in Fig. 15a, finally, shows the coordinate $y_{\sigma_n=0, x=L/2}(P)$ of a zero normal stress $\sigma_{n_x=L/2}(y) = 0$ in the critical cross-section (x = L/2), in dependence of the imposed load *P*. It is shown that this coordinate is going from infinity to zero, while the load *P* increases from zero towards P_{cr} .

Conclusion

The design issue of stability for load-bearing members, in general, is a critical step of the overall design process and should be properly addressed. In this study, the attention was focused on the proposition and validation of a theoretical formulation for the technical stability analysis of slender columns in compression. The theoretical study was thus addressed with the support of case-study calculation examples carried out on slender columns made of steel.

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Ethics

The content of this paper is the changed preprint (Murawski, 2019) according to the issuer requirements and the reviewers and own author's corrections.

Conflicts of Interest

The author declares no conflict of interest.

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